

Global Structure of the ISM in galactic disks and in the central region of galaxies

WADA Keiichi

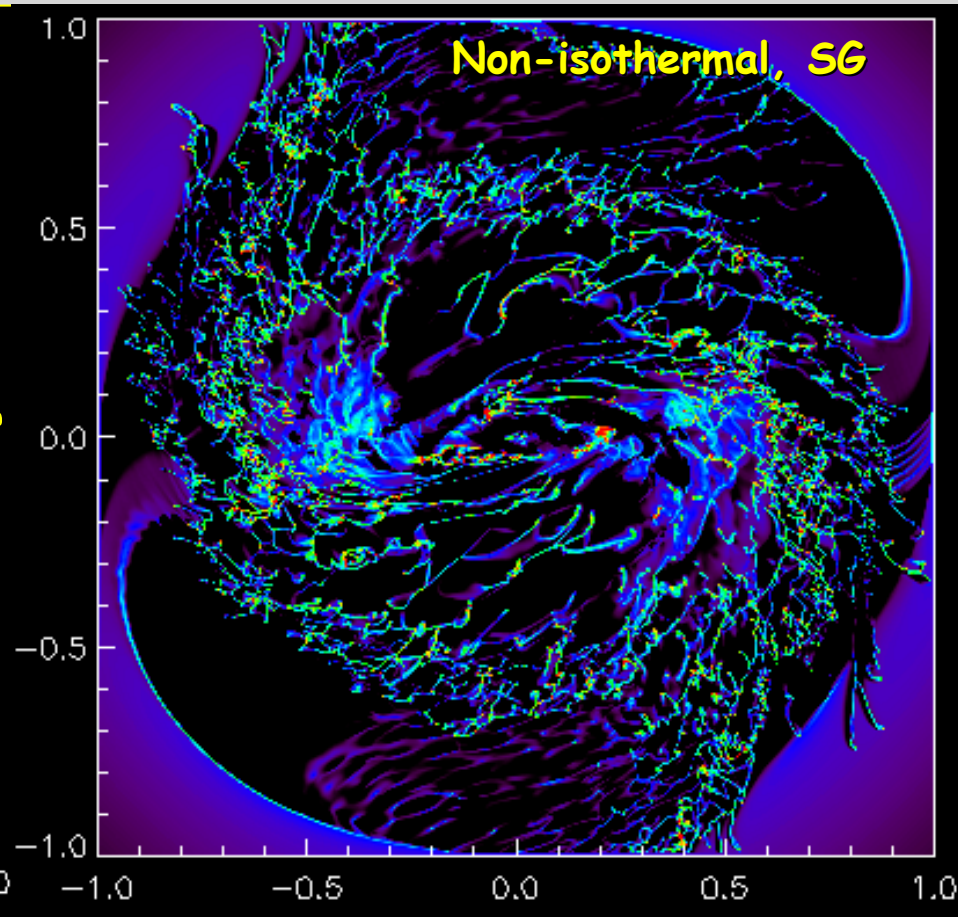
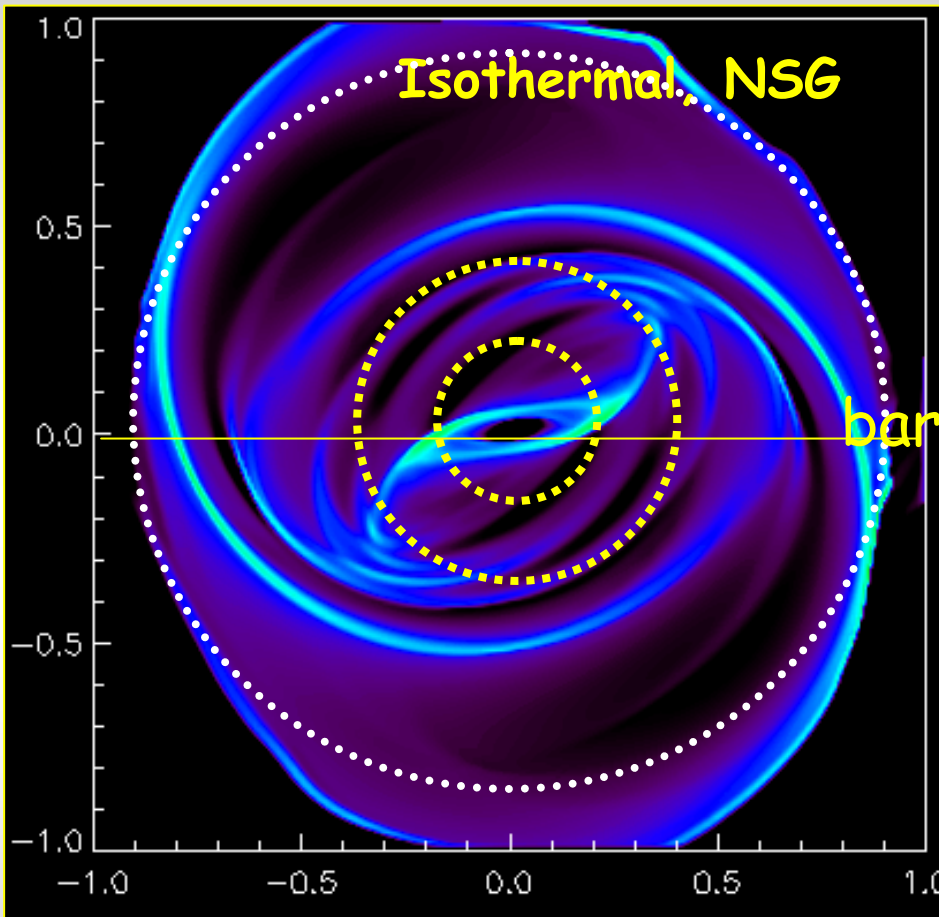
(National Astronomical Observatory of Japan)

Methods & techniques

2D/3D Hydrodynamic Simulations,
Euler method with a uniform Cartesian grid

- * Poisson equation for self-gravity
- * Energy equation with heating/cooling processes
- * Optimization for high-performance computers

Gas dynamics in a bar-potential



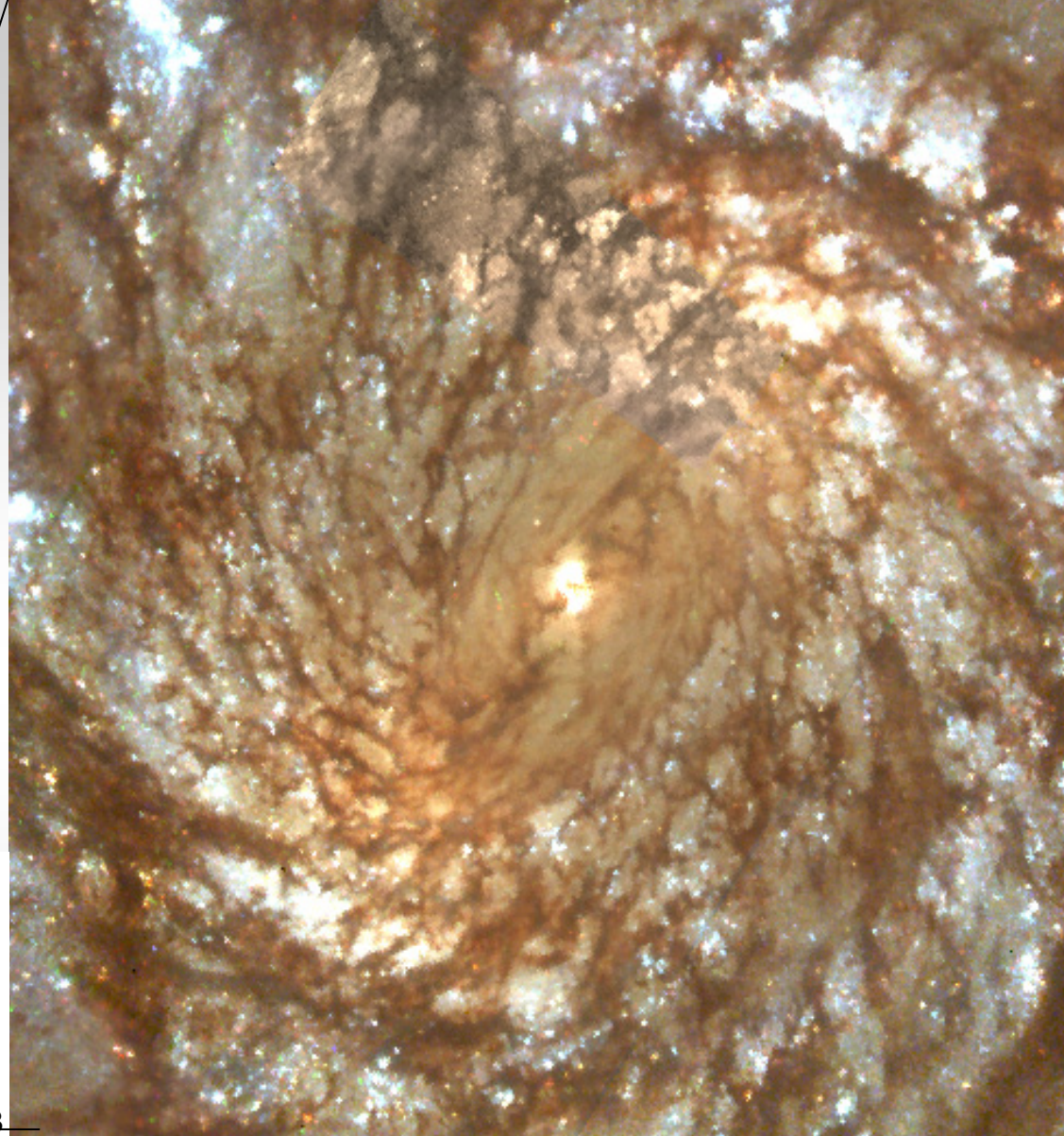
IILR=200pc, OILR=400pc, CR= 900pc

Taking into account relevant physics is essential, but....

Central region of M51 (2kpc x 2kpc)

The ISM is very
inhomogeneous on
~10 pc scale.

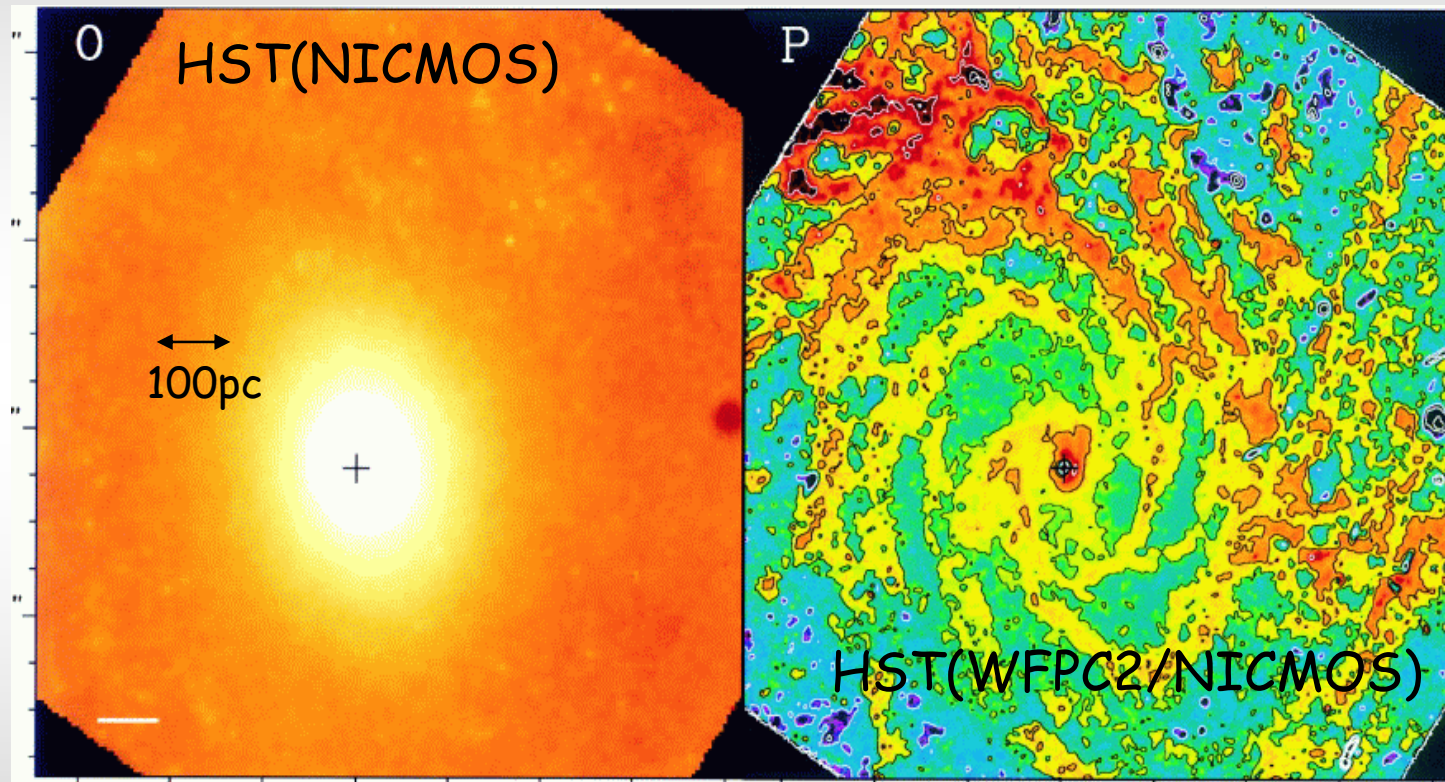
Many filaments, and "mesh"
of the dust lanes



HST obs. (Scoville et. al)

Central region of Seyfert NGC 3982

multi-armed dust lanes



Regan & Mulchaney (1999)

Origin of the structures: Instabilities?

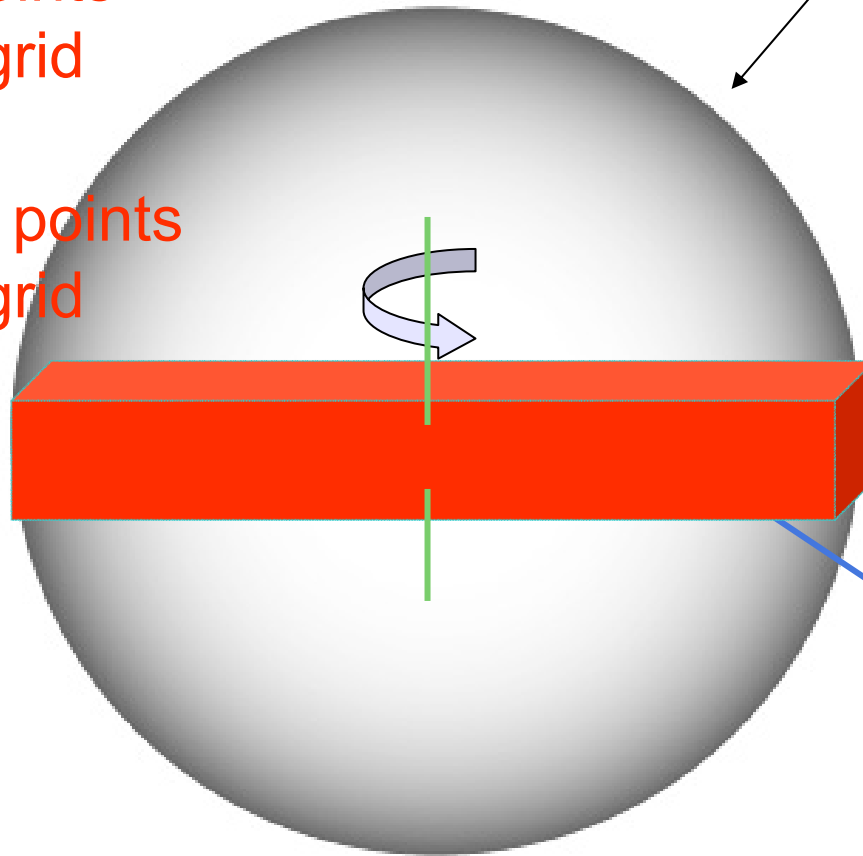
- Filaments/clumps/clouds, holes, voids/holes
- Multi-phase: Cold gas, Warm gas, Hot gas
- Turbulent + rotation
- Is it possible to be quasi-stable on a galactic scale, but unstable on a local scale?
- Can we reproduce these features by numerical simulations from the first principle (with minimal assumptions)?

2D/3-D Hydrodynamics of a gas disk in a spherical galactic potential

2D: 2048^2 grid points
Resolution: 1pc/grid

3D: $512^2 \times 64$ grid points
Resolution: 5pc/grid

Stellar/DM potential: fixed



Thin gas disk
 $\sim 10^{8-9} M_{\text{sun}}$
 $R = 1.28 \text{ kpc}$

Evolution/Structure of the massive gas disk

-- 2-D/3-D Hydrodynamic **Global** Modeling --

Initial conditions and input physics:

- Rotationally supported, uniform disk in a **fixed spherical potential** (stars +DM)
- **Self-gravity** of the gas
- A **cooling** function ($10 < T < 10^8$ K) is assumed.
- Heating sources: uniform UV & supernova explosions
 - Evolution of SNR is directly followed w/ sub-pc resolution

Basic Equations

Conservation of mass, momentum, & energy, Poisson eq.

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{v}) = 0, \quad (1)$$

$$\frac{\partial \mathbf{v}}{\partial t} + (\mathbf{v} \cdot \nabla) \mathbf{v} + \frac{\nabla p}{\rho} + \nabla \Phi_{\text{ext}} + \nabla \Phi_{\text{sg}} = 0, \quad (2)$$

$$\frac{\partial E}{\partial t} + \frac{1}{\rho} \nabla \cdot [(\rho E + p) \mathbf{v}] = \Gamma_{\text{UV}} - \rho \Lambda(T_g), \quad (3)$$

$$\nabla^2 \Phi_{\text{sg}} = 4\pi G \rho, \quad (4)$$

Methods: AUSM w/ uniform grid: 256³, 256³ x 128, 512² x 32; 2048² - 4096²)
+ Poisson eq. Solver(FFT)

CPU time: ~ 10-200 hours/run

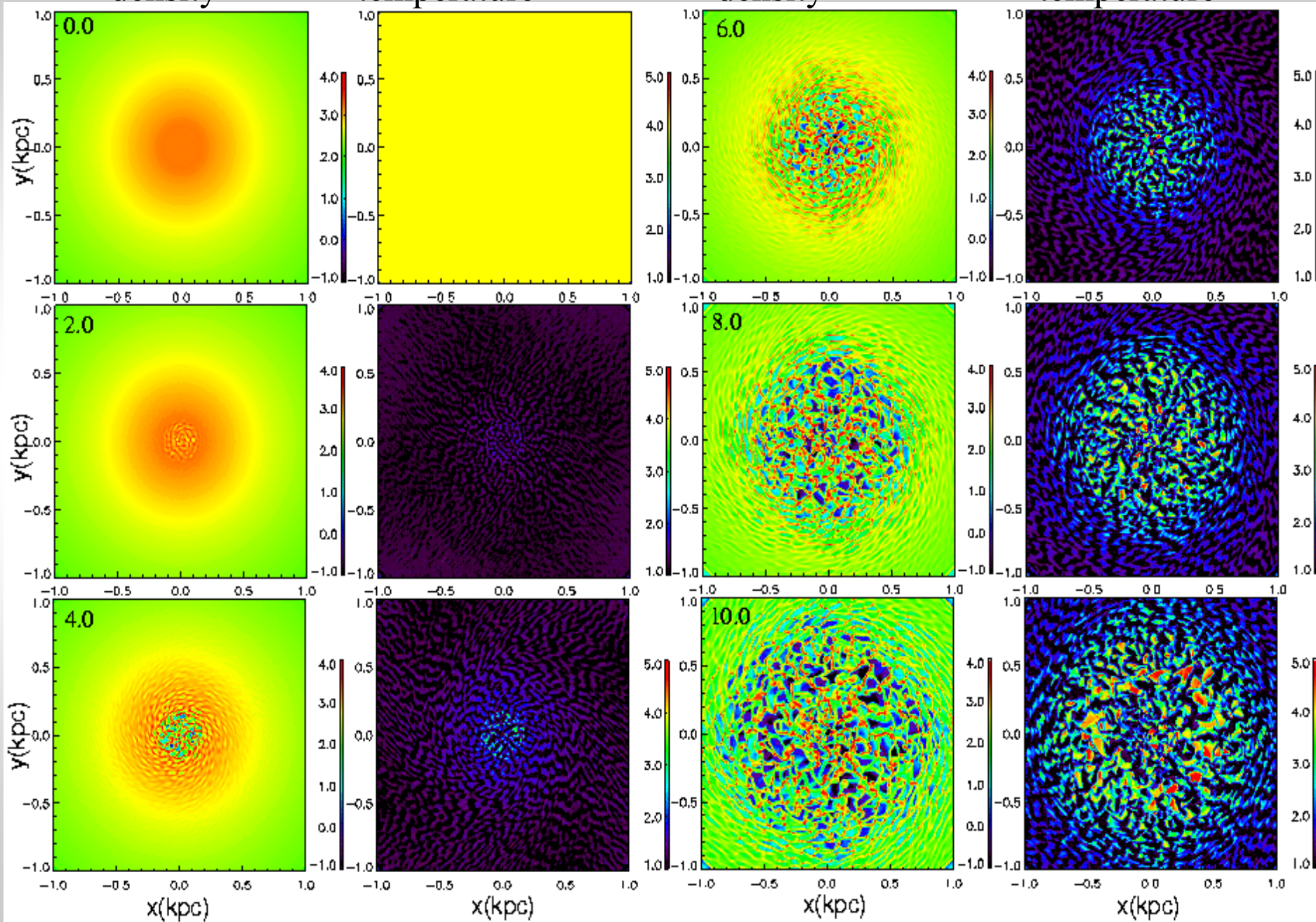
Time evolution of the disk

density

temperature

density

temperature

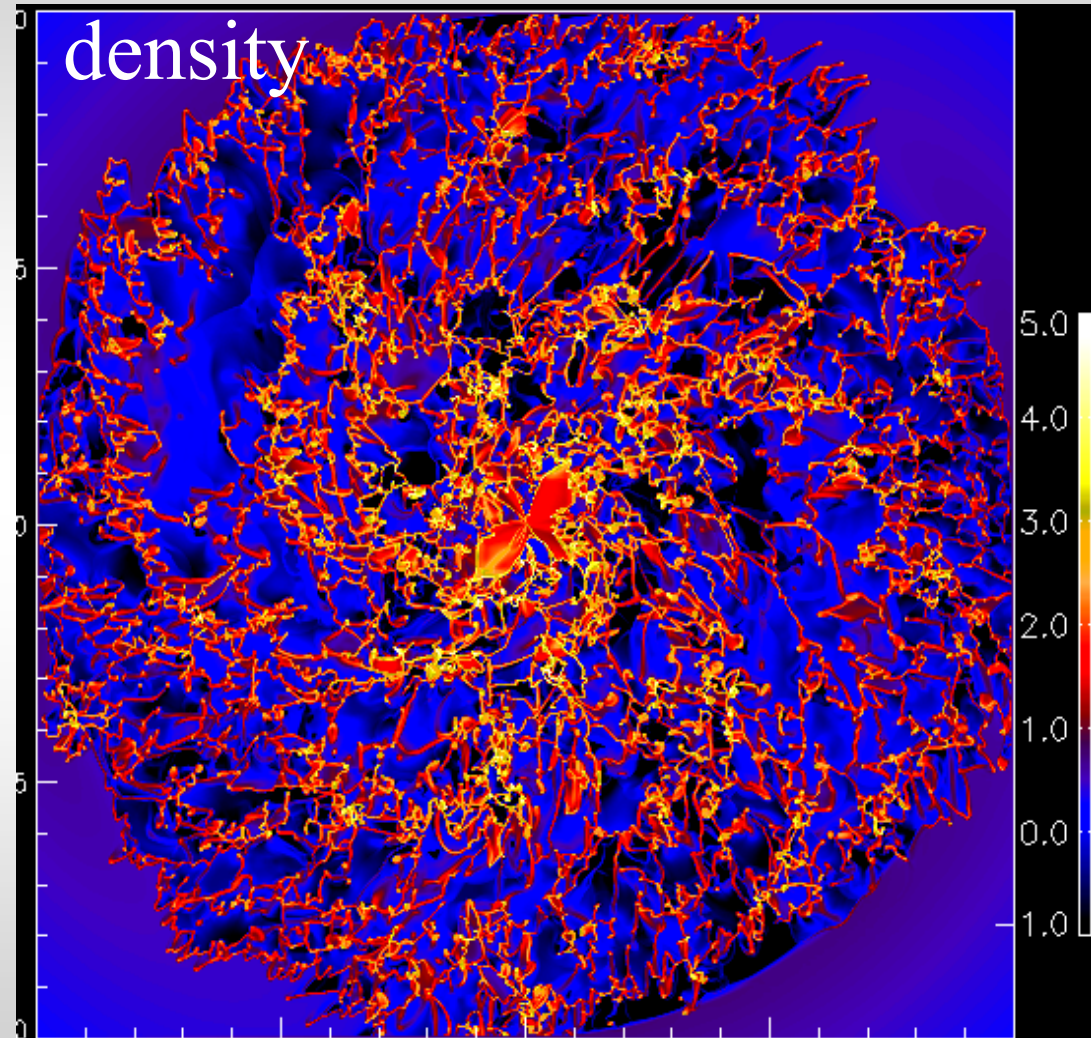


"tangled-web" structure of the ISM in a galaxy

No energy input from supernovae

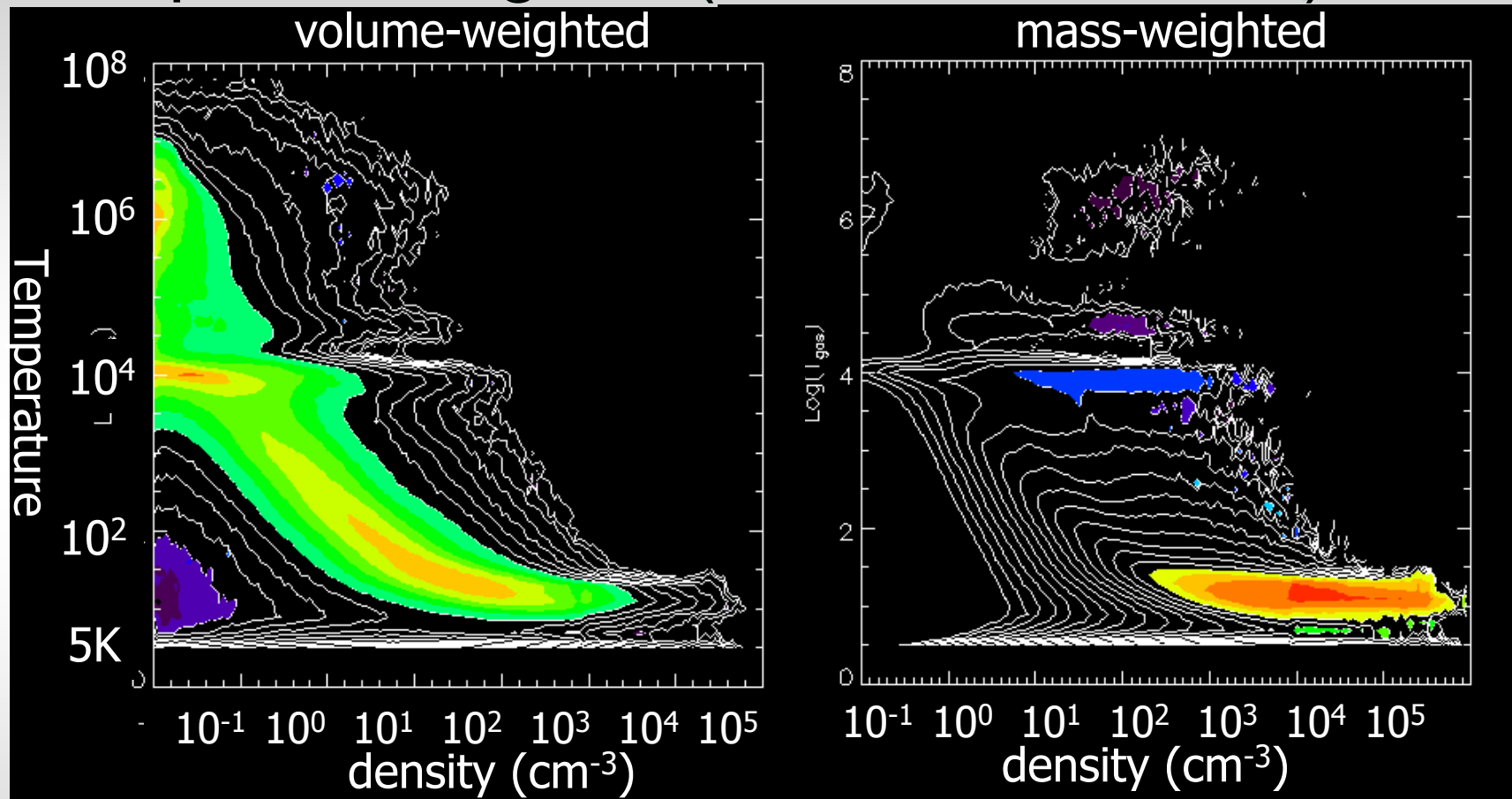
2048² grids, 0.98pc/grid

- *Clumps/filaments*
 - *high density, low T ($<100\text{K}$)*
 - *"GMC" = complex of clouds & filaments*
- *Cavities/Holes*
 - *low density, high T*
 - *Shock heated gas ($\sim 10^5\text{K}$)*



Wada & Norman, ApJ 516, L13 (1999)
ApJ 546, 172 (2001)

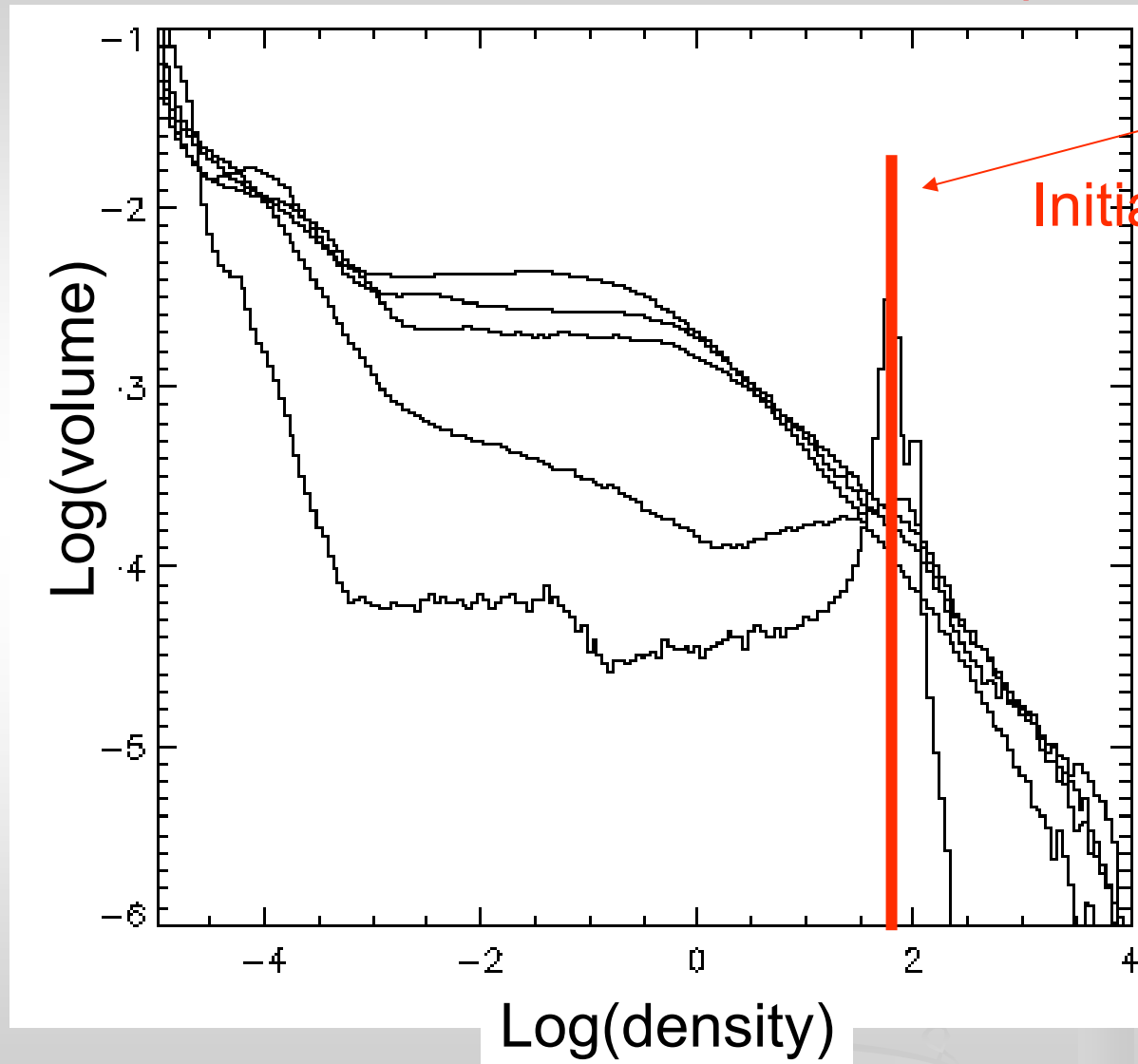
physical condition of the ISM phase-diagram (with SN feedback)



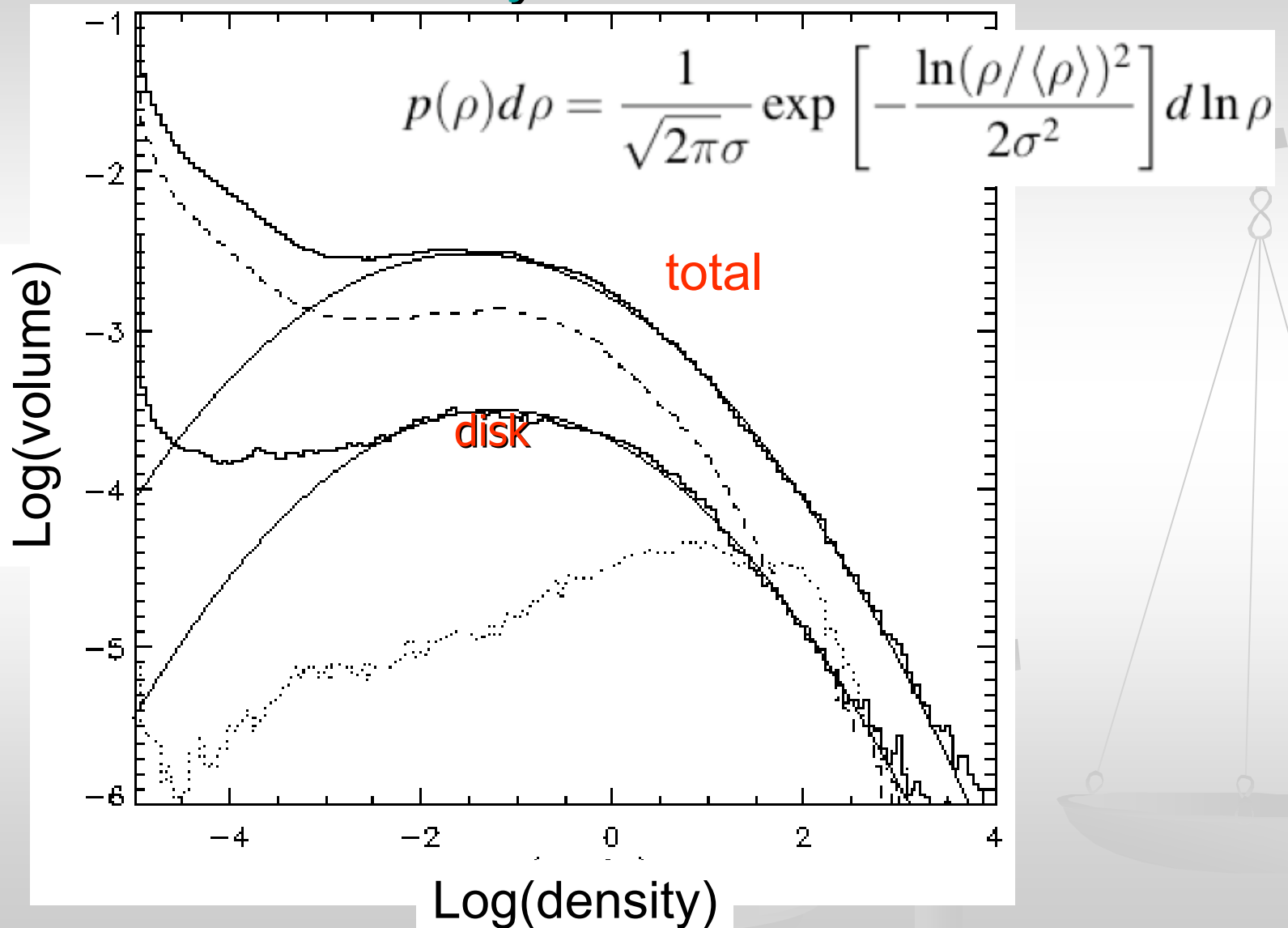
- Density ranges over four orders of magnitude for a given T_{gas} .
- $\Rightarrow$ Pressure is distributed in a wide range
- $\Rightarrow$ Different phases are not necessary in a pressure equilibrium

Evolution of PDF in 3-D galactic disk

Initial condition is lost within a few dynamical times.



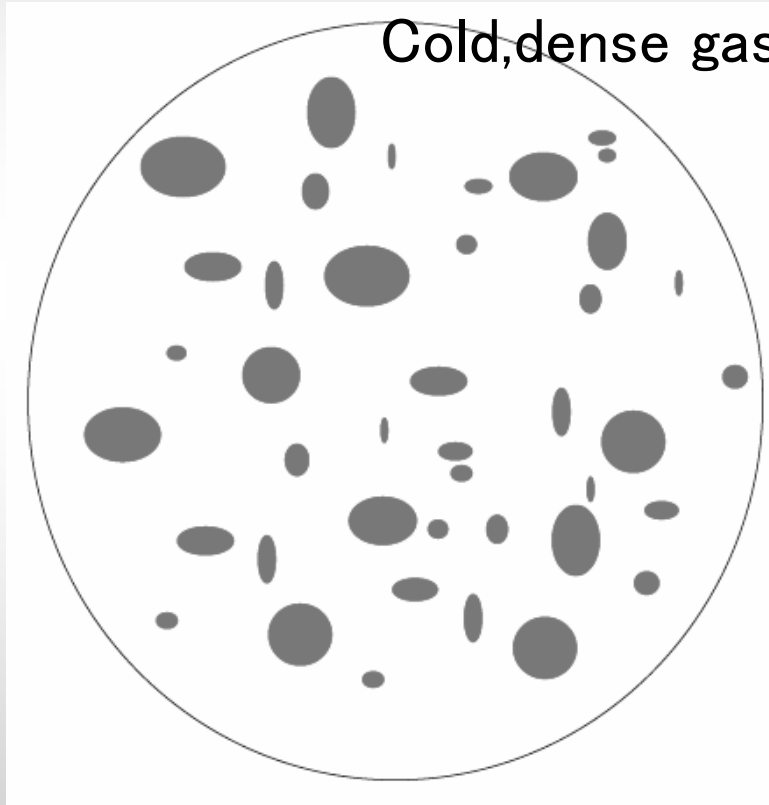
Volume-weighted density PDF in a steady-state



ISM in a galactic disk is NOT discrete phases.

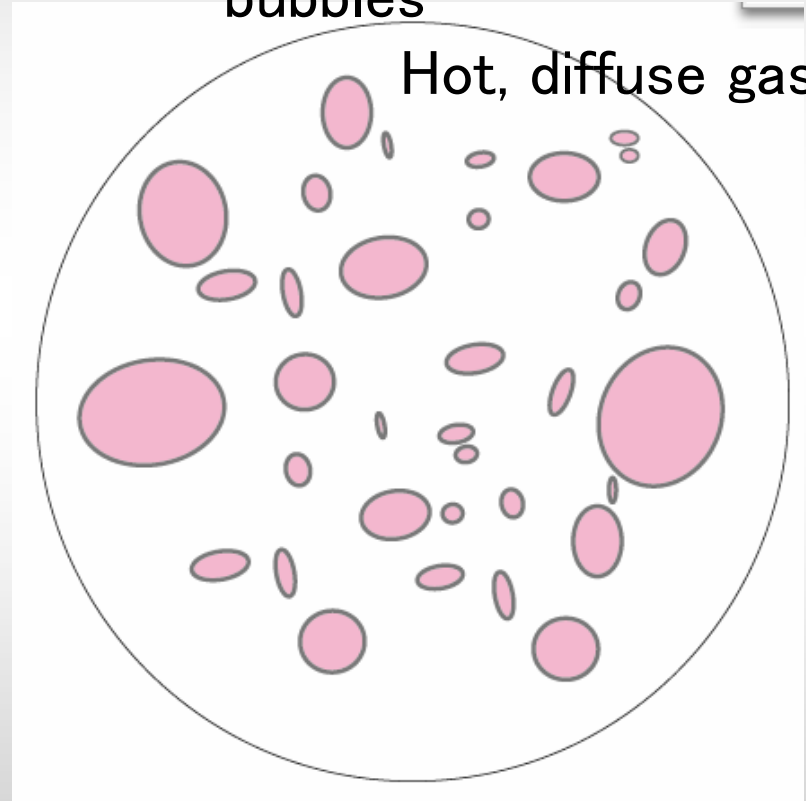
Cloudlets

Cold, dense gas



bubbles

Hot, diffuse gas



High density gas is not independent of lower density gases.

Universal PDF of the ISM in galaxies

Log-Normal part: Highly inhomogeneous.

Higher density gases occupy smaller volumes.

Structures of dense gases are **not independent** of lower density gases.

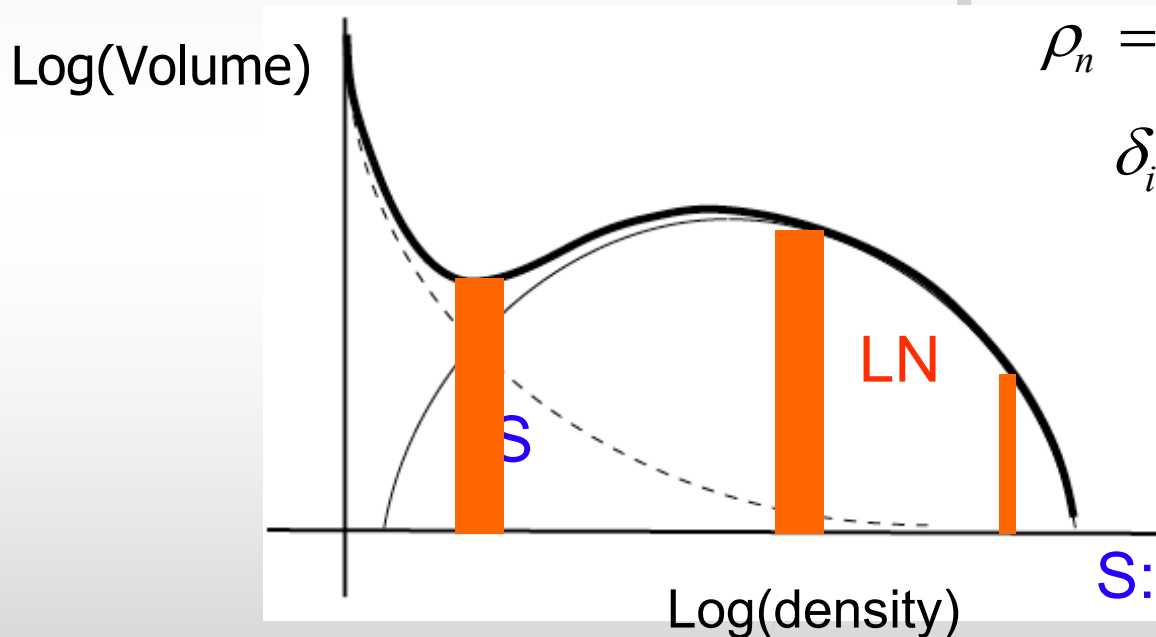
Shock, collapse, etc $\rightarrow \rho_i = \delta_{i-1} \rho_{i-1}$

$$\rho_n = \delta_n \delta_{n-1} \dots \delta_0 \rho_0$$

δ_i : independent events

$$\ln(\rho_n) \rightarrow N(\mu, \sigma^2) \\ n \rightarrow \infty$$

Central limit theorem



S: Smoothed part

Globally stable disk $\rightarrow$ PDF does not evolve

A 2-D model with energy feedback from supernovae

1 kpc

density



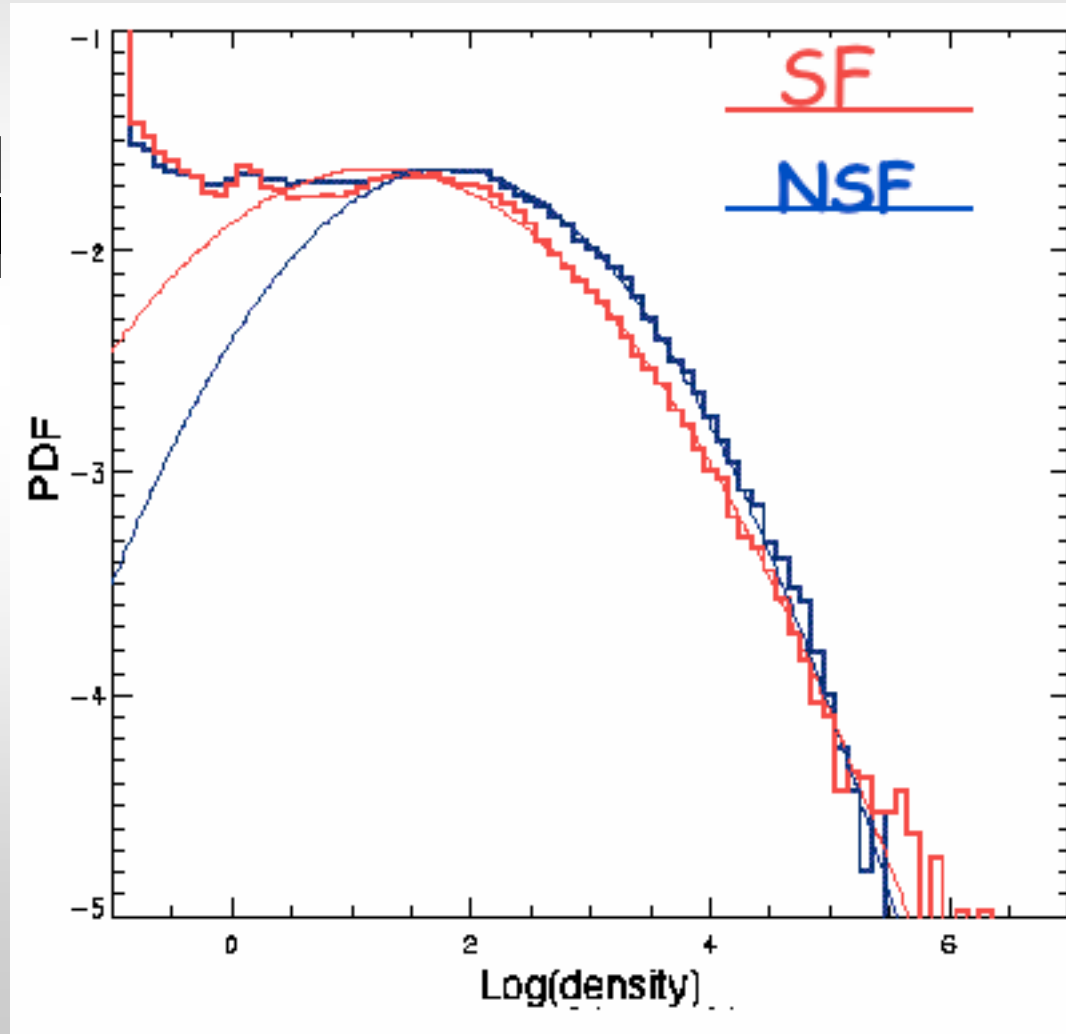
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temperature
Red: $T > 10^6$ K

Effect of stellar energy feedback on the PDF

Log-Normal PDF is robust

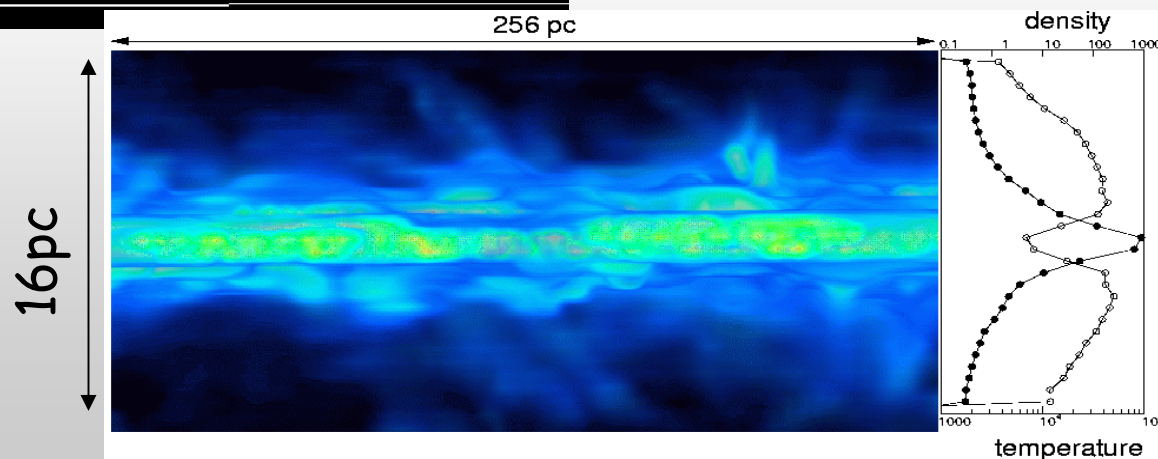
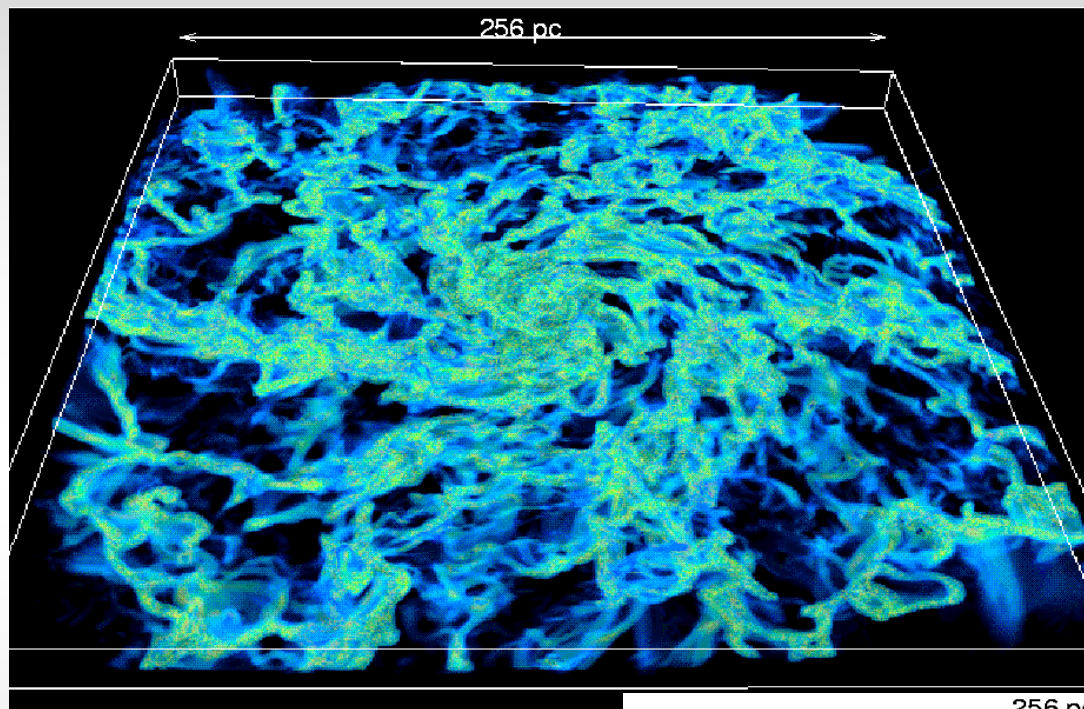
$$PDF(\rho) \propto \text{Exp}\left[\frac{-(\ln(\rho) - \ln(\rho_0))^2}{2\sigma^2}\right]$$

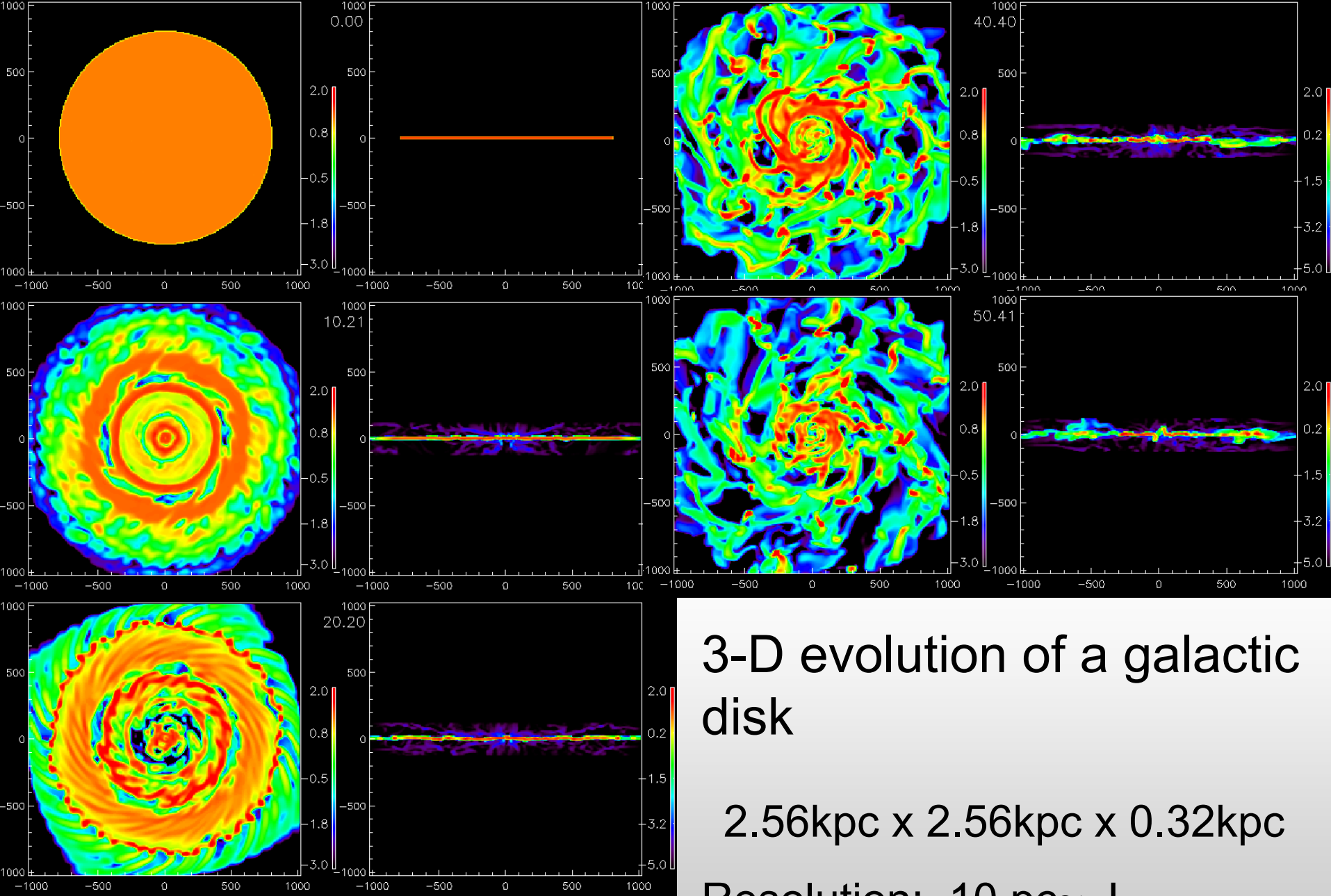


3D structure of a nuclear massive disk

512² x 32 grid points (0.5 pc/grid)

Wada (2001) ApJ 559, L41





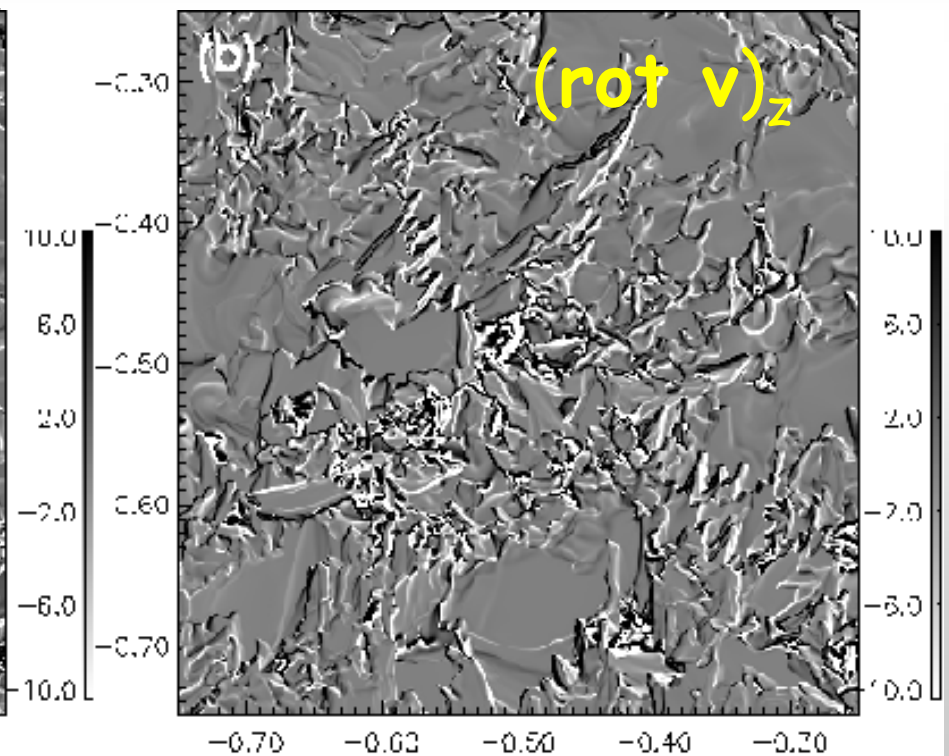
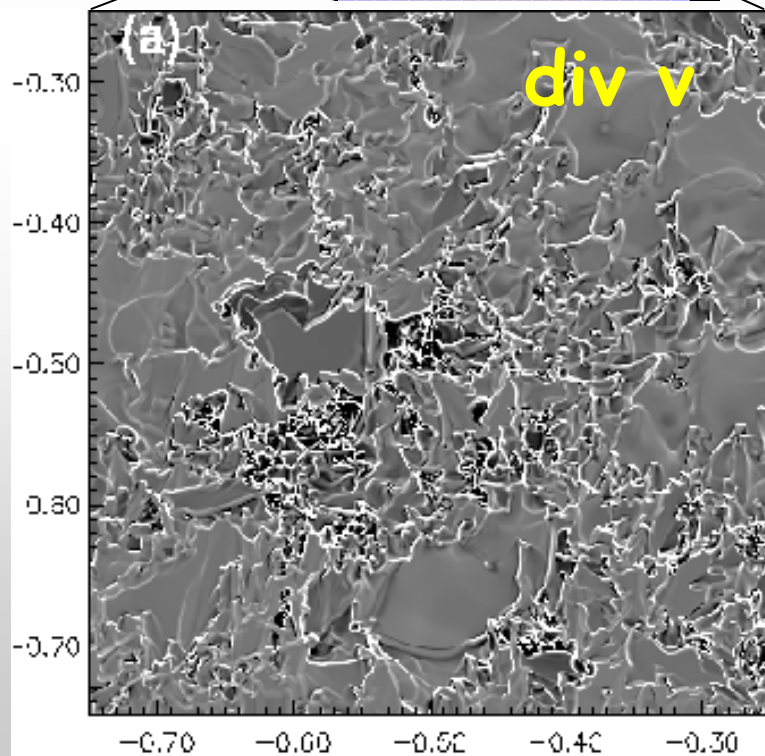
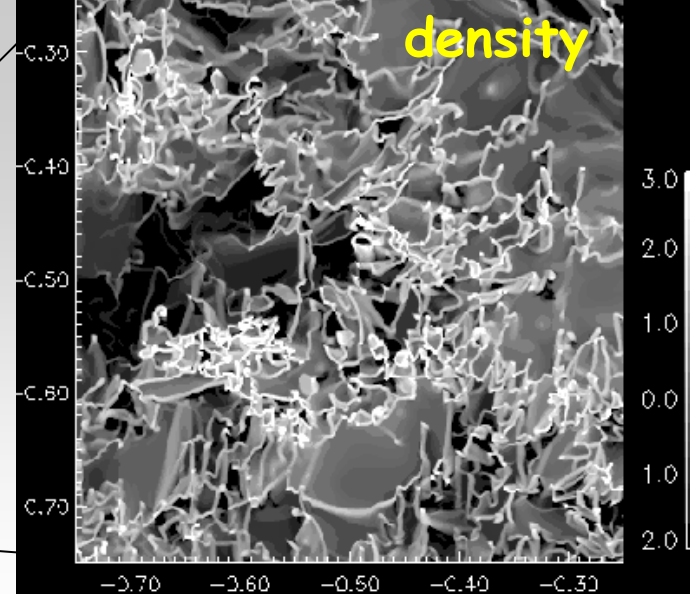
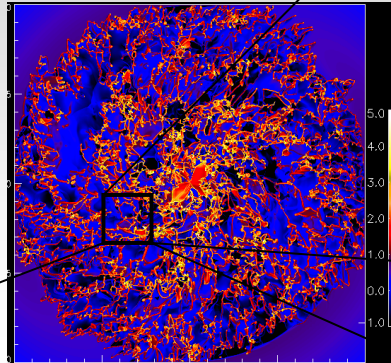
3-D evolution of a galactic disk

2.56kpc x 2.56kpc x 0.32kpc

Resolution: $10 \text{ pc} \sim L_{\text{jeans, min}}$

Turbulent structure of the 'tangled web'

No supernovae

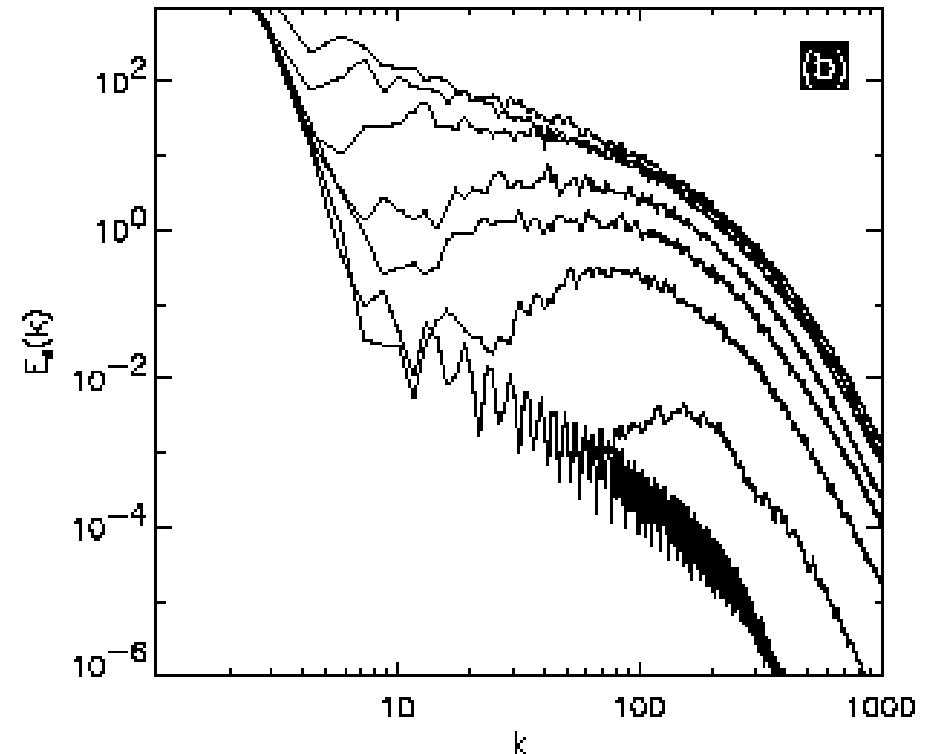
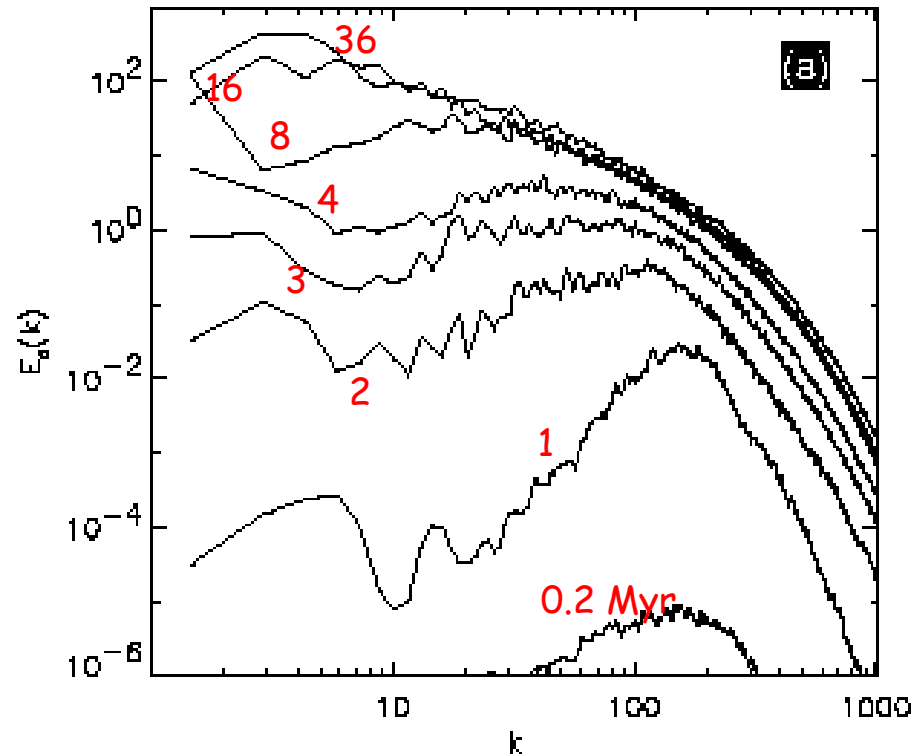


Evolution of Energy spectrum (no SN, 2-D model)

The spectrum attains a power-law in $\sim 20\text{Myr}$

Compressible (rotational free)

Incompressible (divergence free)



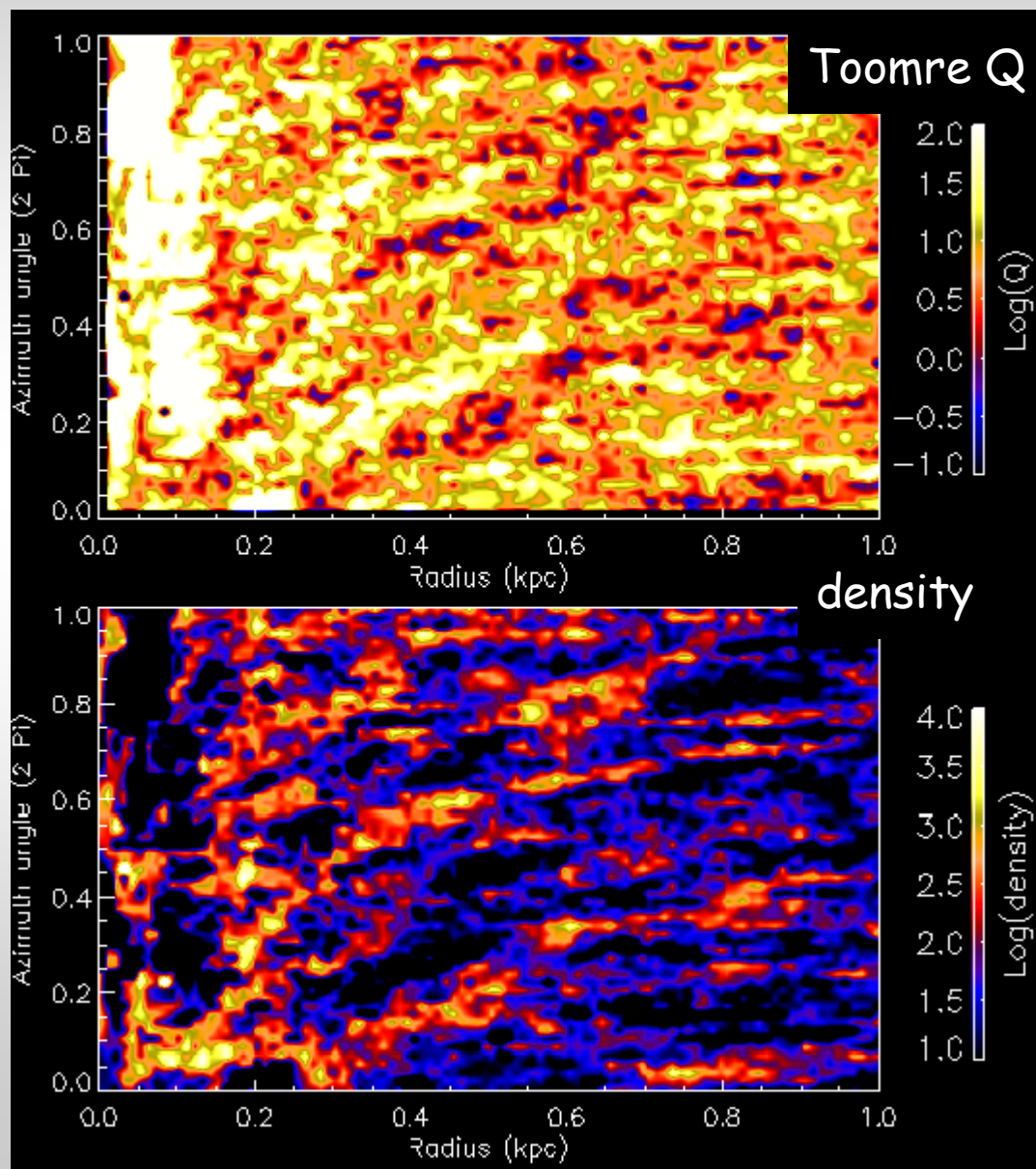
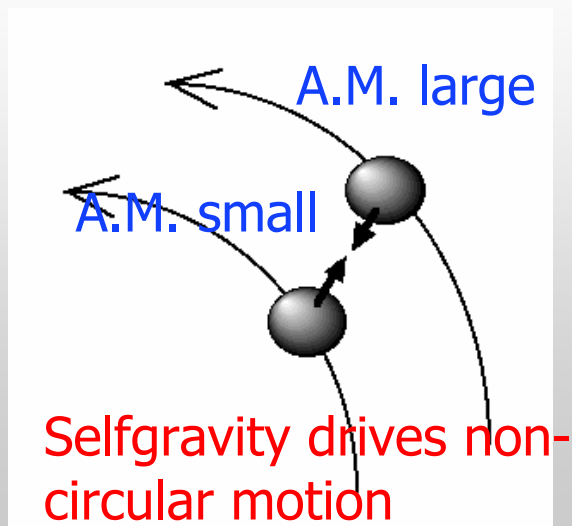
There is **no explicit energy input**.

Turbulence on a galactic scale is self-regulated without supernovae.

Rotational component: galactic rotation dominates turbulent energy on large-scale

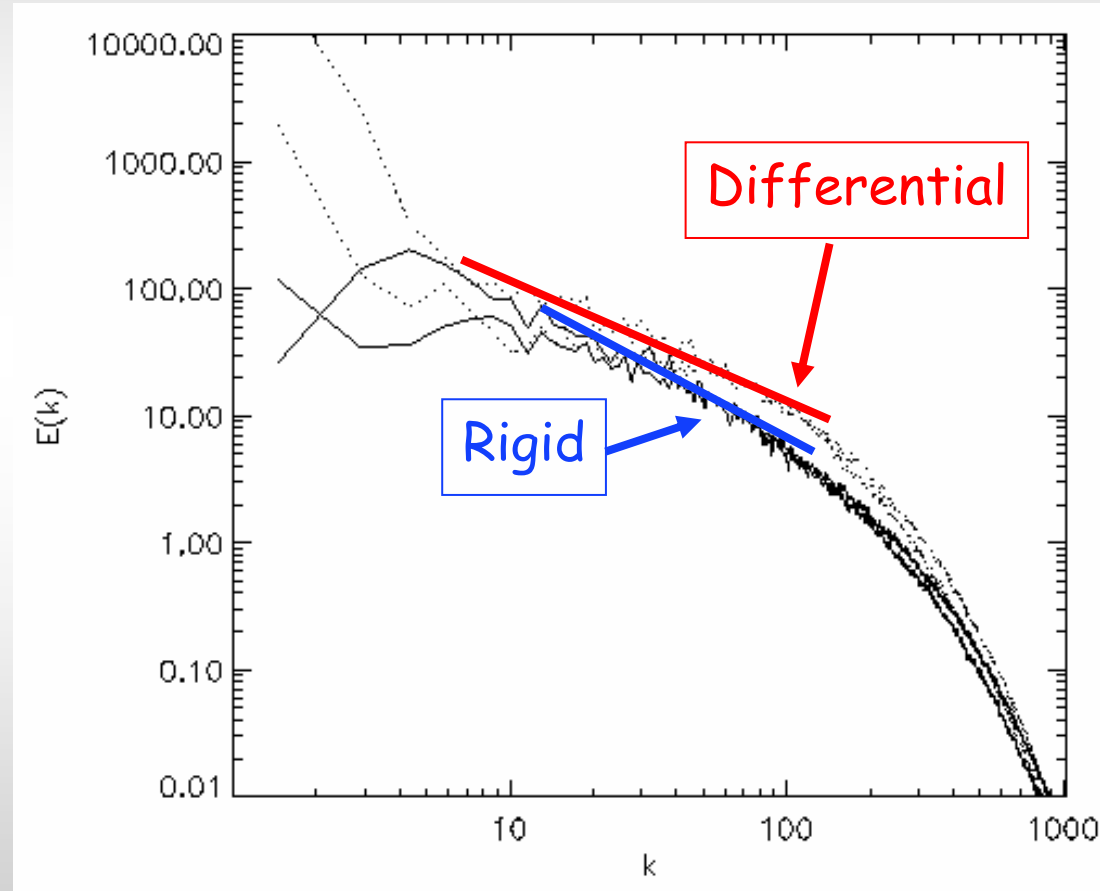
Distribution of stable/unstable regions

- Stable and unstable regions patchily exist.
- median $Q \sim 5$, but large dispersion
- Energy source of the turbulence = local gravitational instability + galactic rotation

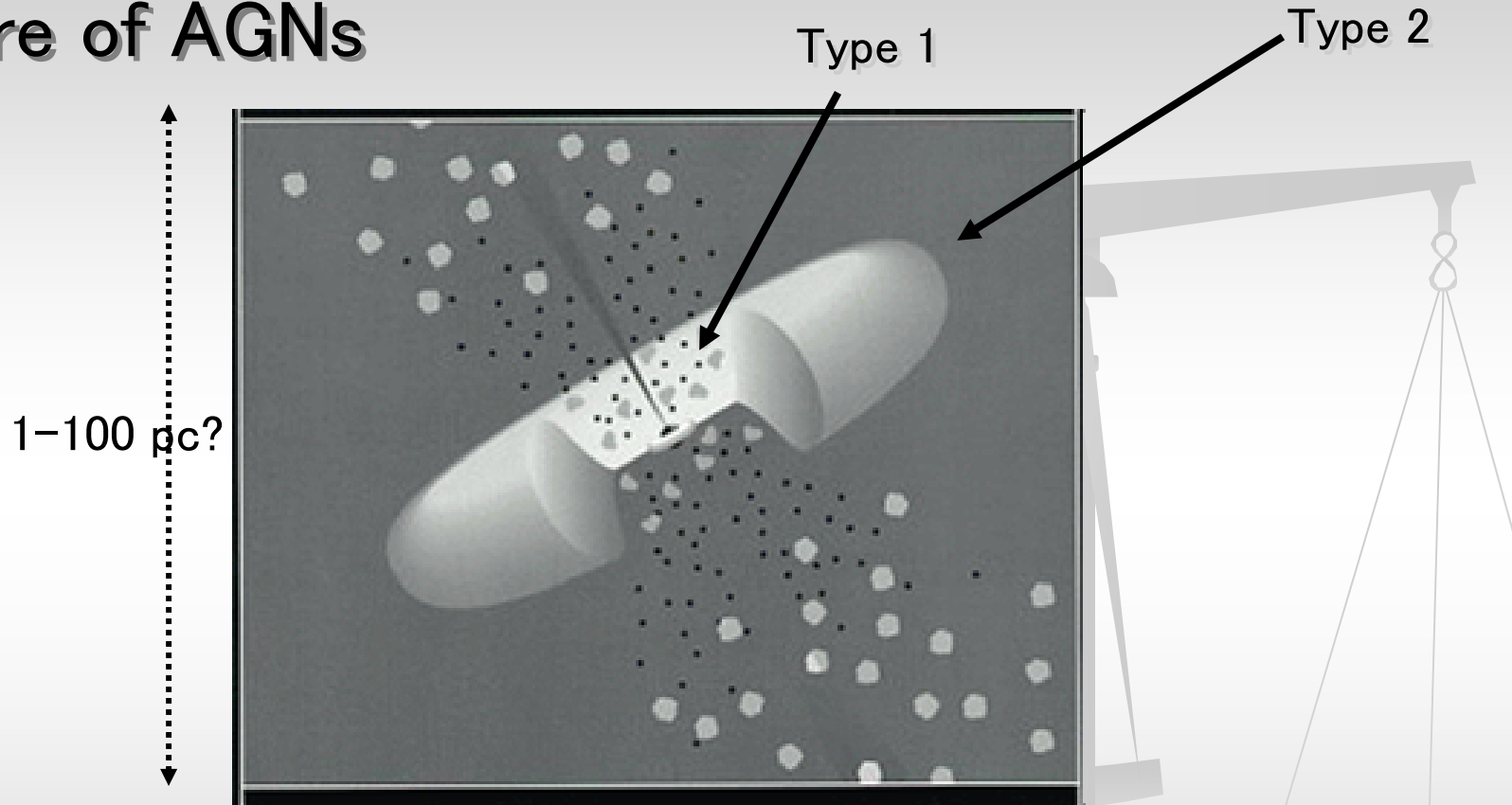


Note: differential rotation is not essential to maintain the turbulence

- Turbulence is also developed in a rigid rotation disk.
- Energy spectrum is slightly steeper in rigid rotation.



Structure of AGNs



Super massive BH ($10^6\text{--}10^8 M_{\text{sun}}$) + accretion disk (AU scale) + BLR + NRL +
Obscuring molecular torus (1-100 pc) + Jet + ENLR

Energy source: mass accretion ($0.01\text{--}1 M_{\text{sun}} \text{ yr}^{-1}$)

3-D Hydrodynamics of a massive gas disk around a SMBH (KW & Norman 02; KW & Tomisaka 05)

256³ grid points

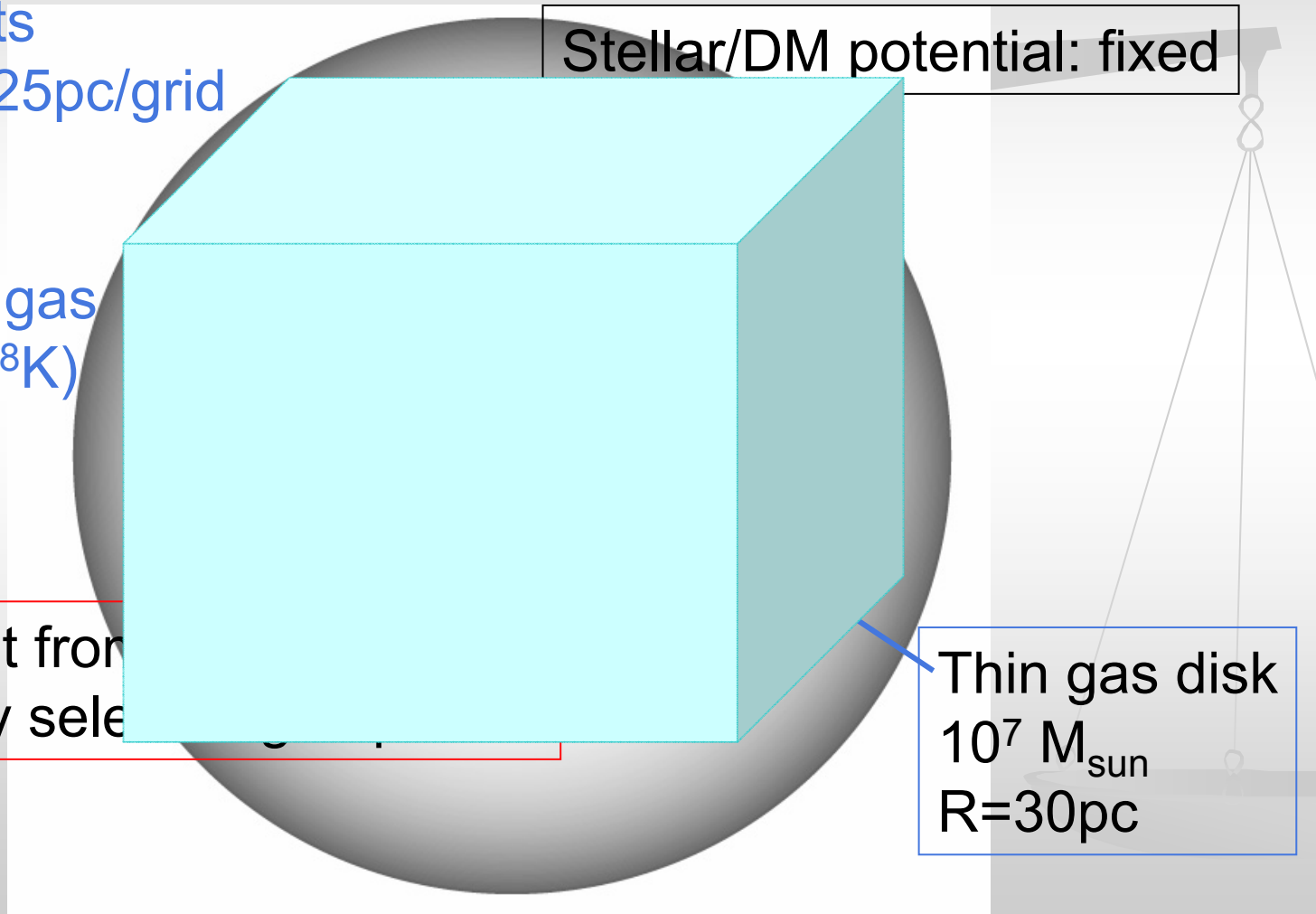
Resolution: 0.25pc/grid

Self-gravity of gas
Cooling(10^{-10} K)

Stellar/DM potential: fixed

Energy input from
at randomly selected

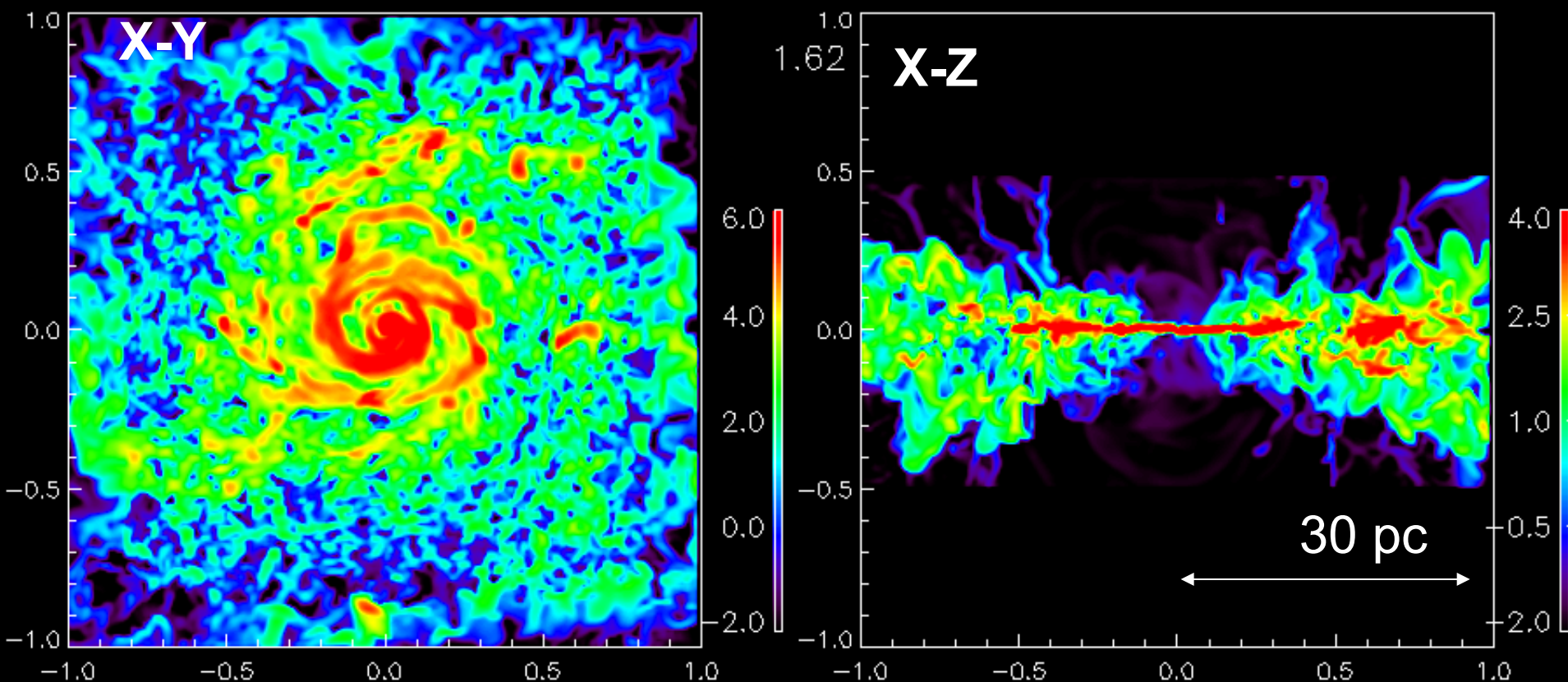
Thin gas disk
 $10^7 M_{\text{sun}}$
 $R=30\text{pc}$



Starburst driven "torus" around SMBH

density

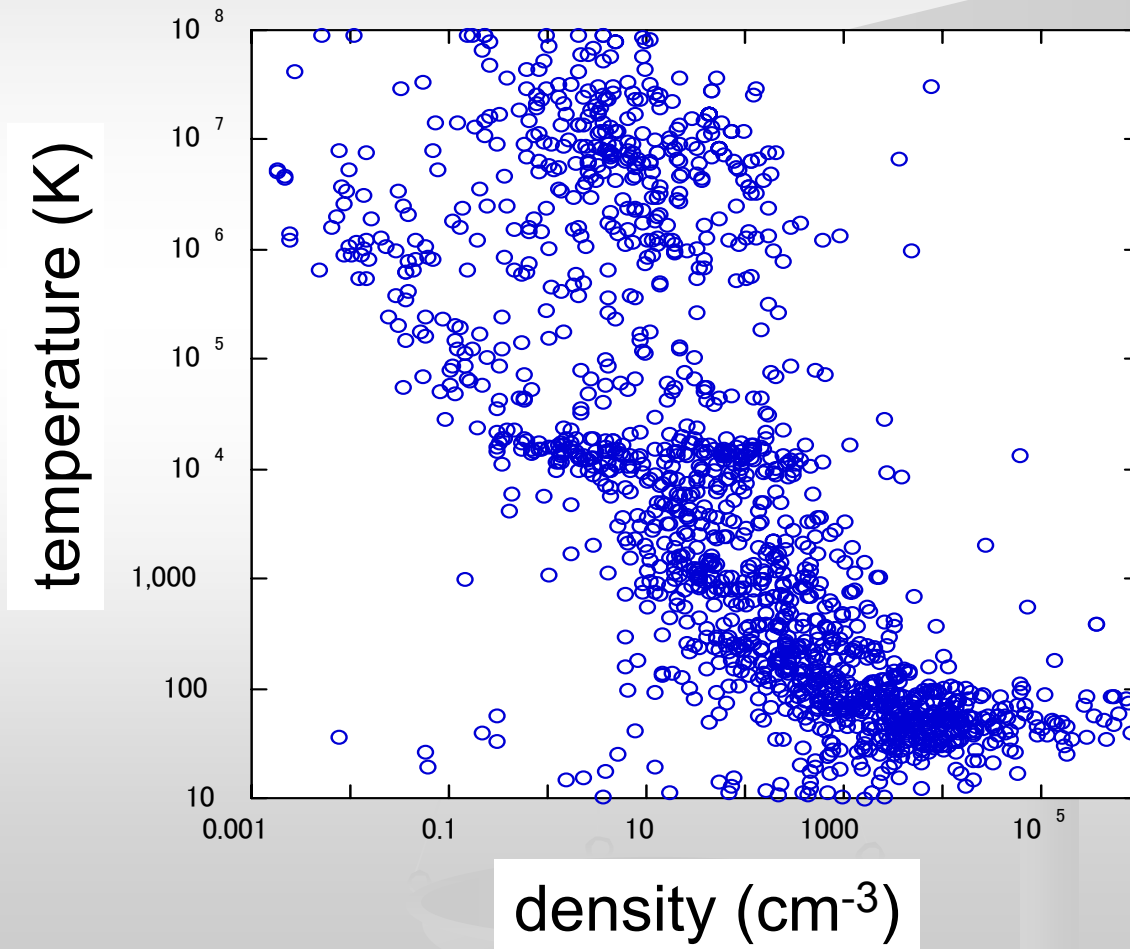
KW & Norman '02



- ISM is a flared disk with clumps and filaments.
- Scale height is determined by energy balance between turbulent dissipation and SN heating in a gravitational potential.

Density, temperature in the “torus”

$T \sim 50\text{-}10^7$ K, $n \sim 0.1\text{-}10^5$ cm $^{-3}$



Obscuring “torus” around a supermassive BH with nuclear starburst

$256^2 \times 128$, uniform grid, 0.25pc/grid

Radiative cooling ($5-10^8$ K), SN feedback, selfgravity

$M_{\text{BH}} = 10^8 M_{\text{sun}}$, $M_{\text{gas}} = 10^7 M_{\text{sun}}$,

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←-----→
64 pc

Geometry of SNe-driven torus

If we assume:

- Internal motion of the torus is turbulent
- turbulent energy dissipation = heating by SNe

$$\frac{\rho_g v_t^2}{\tau_d} = \eta s_{\star} E_s,$$

Scale height $h(r)$



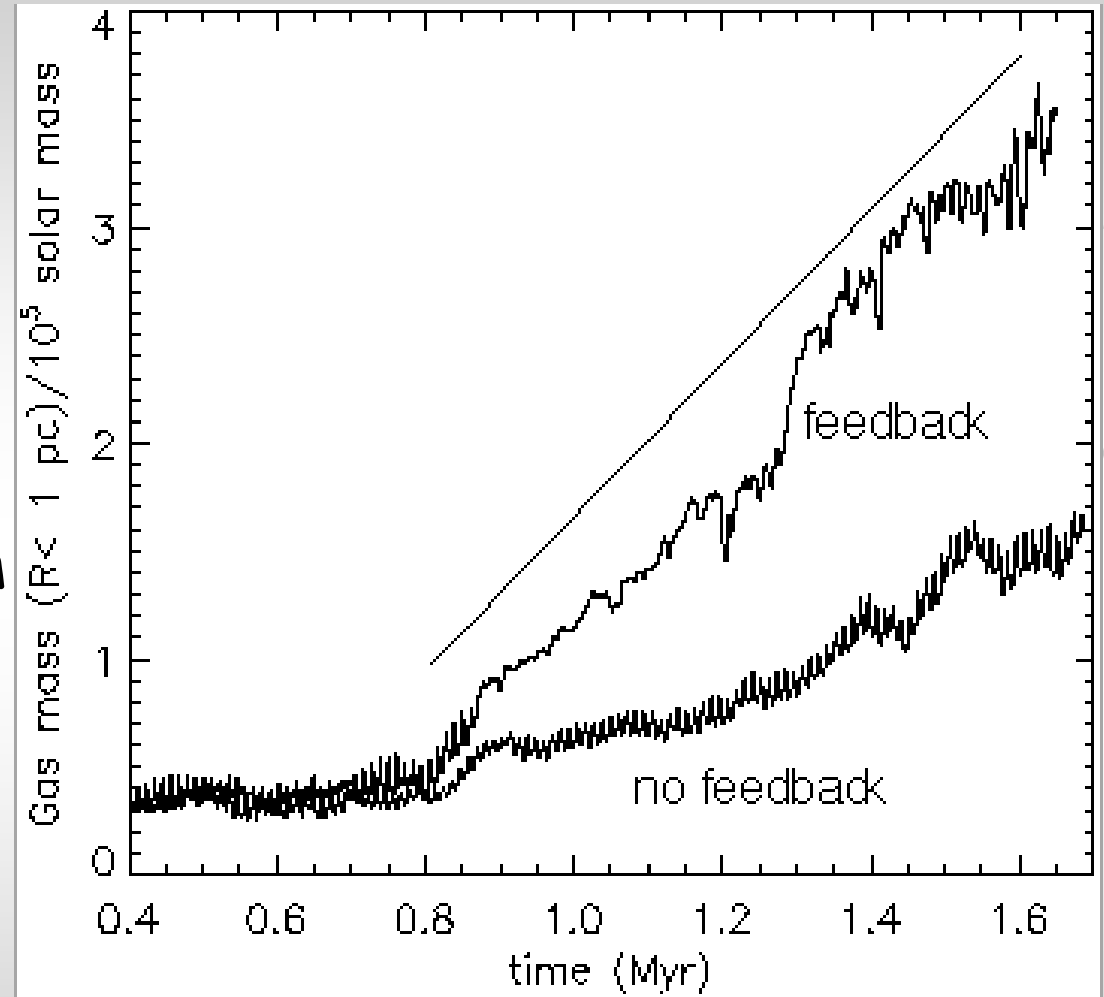
Hydrostatic equilibrium

$$h_1(r) = (\xi \eta \alpha E_s)^{1/2} G^{-1/2} M_{\text{BH}}^{-1/2} r^{3/2}, \quad (\text{BH dominant region})$$

Scale height of the torus
is proportional to $(\text{SFR}/M_{\text{BH}})^{1/2}$

Gas accretion to the nucleus ($R < 1 \text{ pc}$)

- I. average accretion rate $\sim 0.3 \text{ M}_{\text{sun}}/\text{yr}$
- II. accretion rate is enhanced by the starburst.
- III. Angular momentum transport by turbulent viscosity



※ accretion rate is time-dependent ($\tau \sim 10^4\text{-}5 \text{ yr}$)

Summary: Origins of the inhomogeneity of the ISM in galactic disks and in AGNs

1. Gravitational and Thermal instabilities coupled with the galactic rotation
 - Fully developed turbulence is naturally generated in galactic disks.
 - The turbulent motion is maintained by the galactic rotation, and local shear motion.
2. Supernova explosions in a dense gas.
 - SNRs interact with the inhomogeneous ISM, and it causes turbulent motion in the torus.
 - The accretion rate toward the central BH is time dependent.

